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# Testing the dark origin of neutrino masses with oscillation experiments

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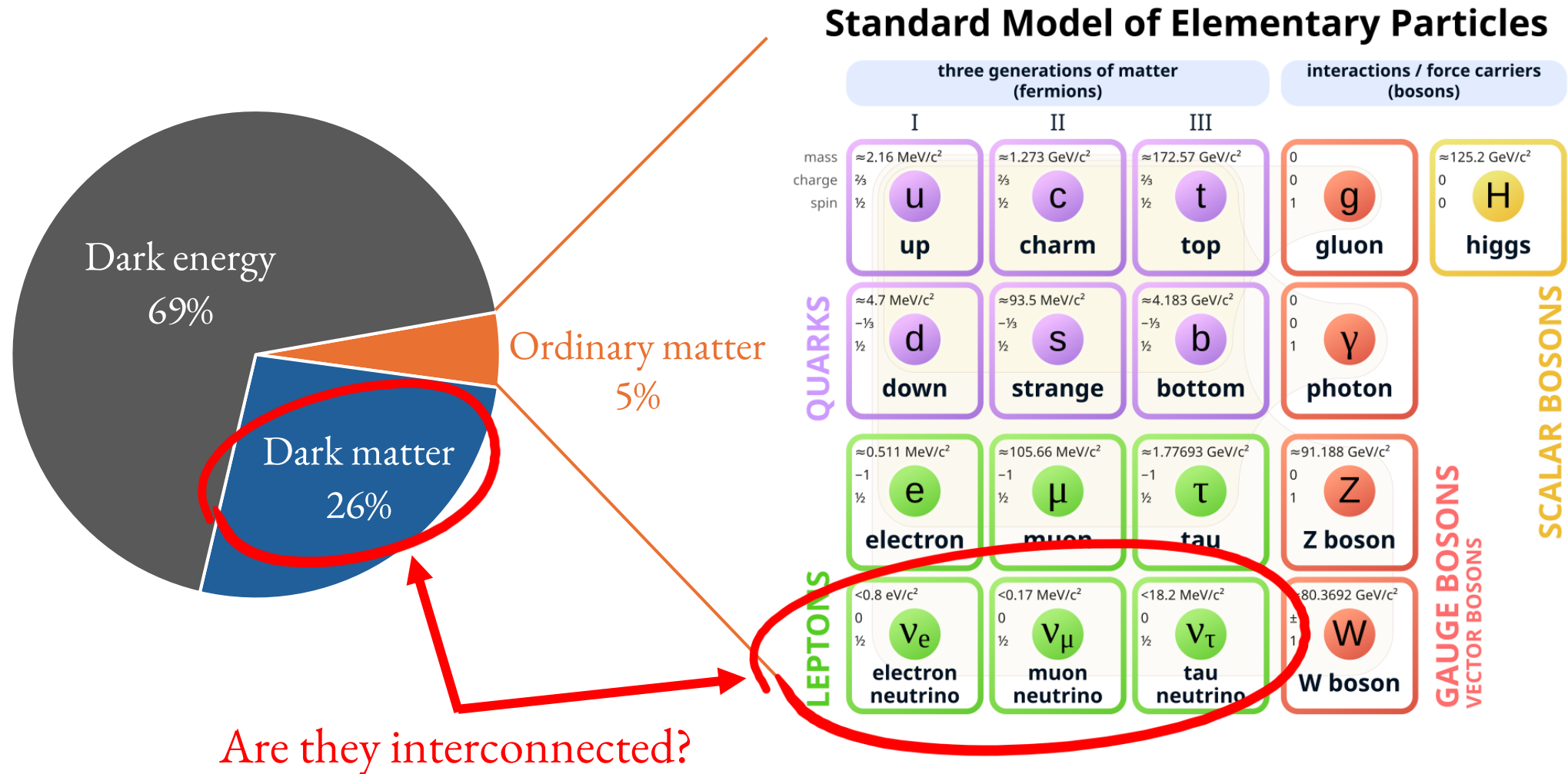


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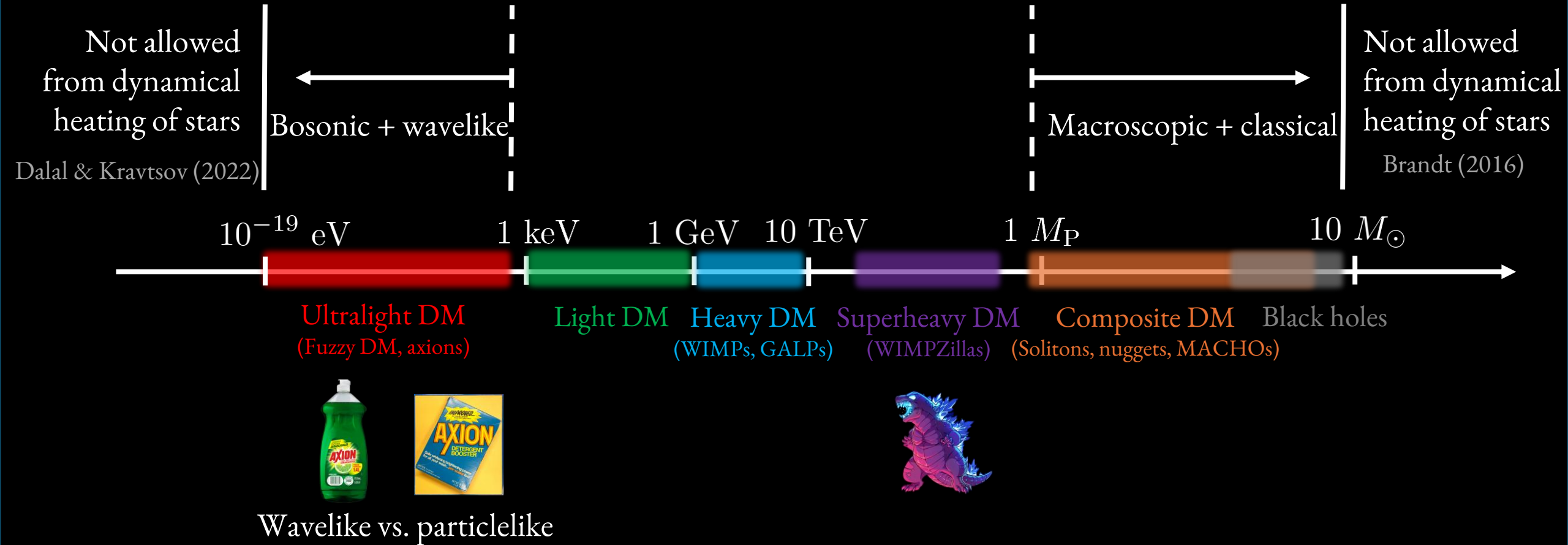


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# Standard model of cosmology and particle physics



# Dark matter mass landscape



# Ultralight dark matter

- Large occupation number → Classical fields

$$n\lambda_{\text{dB}}^3 \sim \left(\frac{40 \text{ eV}}{m}\right)^4 \sim 3 \times 10^{82} \left(\frac{10^{-19} \text{ eV}}{m}\right)^4$$

- Macroscopic/astrophysical scales

$$\lambda_{\text{dB}} \sim 50 \text{ } \mu\text{m} \left(\frac{40 \text{ eV}}{m}\right) \sim 0.6 \text{ pc} \left(\frac{10^{-19} \text{ eV}}{m}\right)$$

- Wave dynamics, rich phenomenology

Suppressed small-scale structure, interference,  
Bose-Einstein condensates, polarization,  
modulations of standard model “constants”, etc.



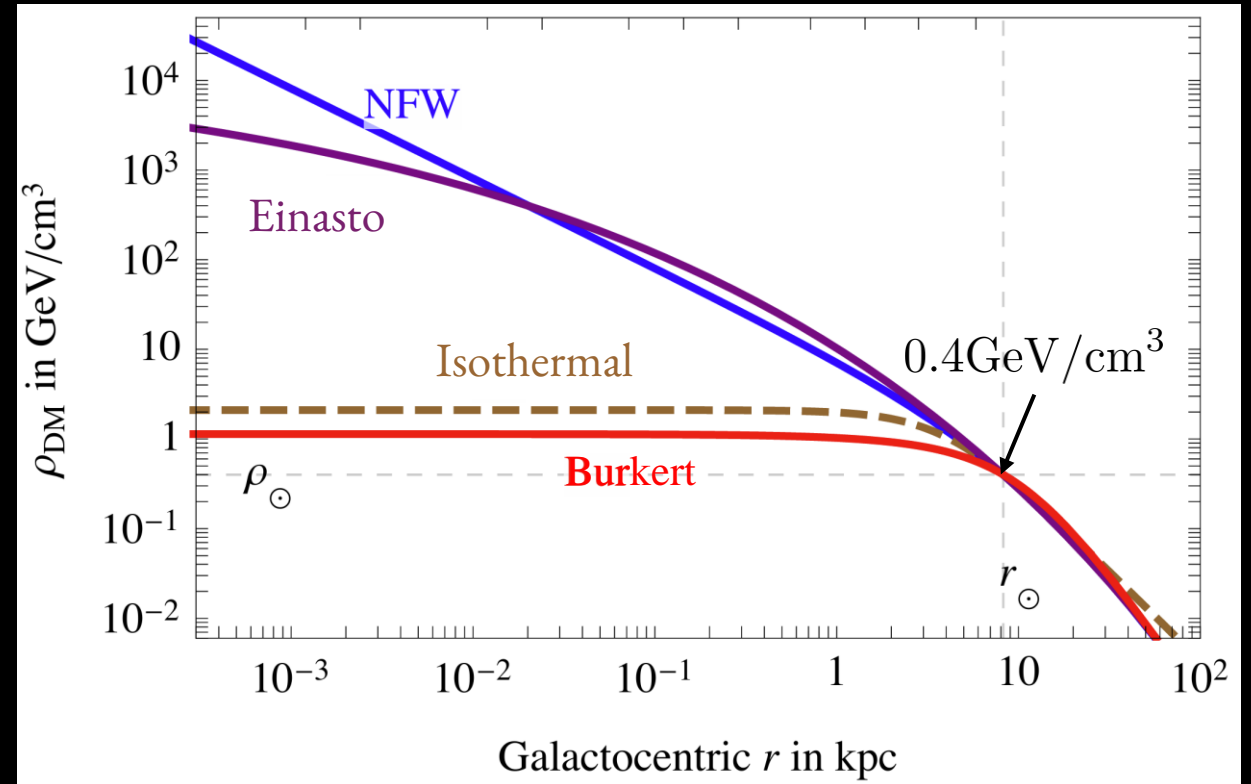
assuming  $\rho \sim 0.4 \text{ GeV}/\text{cm}^{-3}$ ,  $v \sim 200 \text{ km/s}$

# Density profiles and fluctuations

Wave interference, causing  $\gtrsim \mathcal{O}(1)$  fluctuations in local density.

$$\phi(t, \mathbf{x}) \simeq \phi_0(\mathbf{x}) \cos(m_\phi t)$$

$$\left\{ \begin{array}{l} L \ll \lambda_{\text{dB}} : \phi_0(\mathbf{x}) \approx \phi_0 \\ L \gg \lambda_{\text{dB}} : \text{stochastic } \phi_0(\mathbf{x}) \end{array} \right.$$



Cirelli et al. (2024)

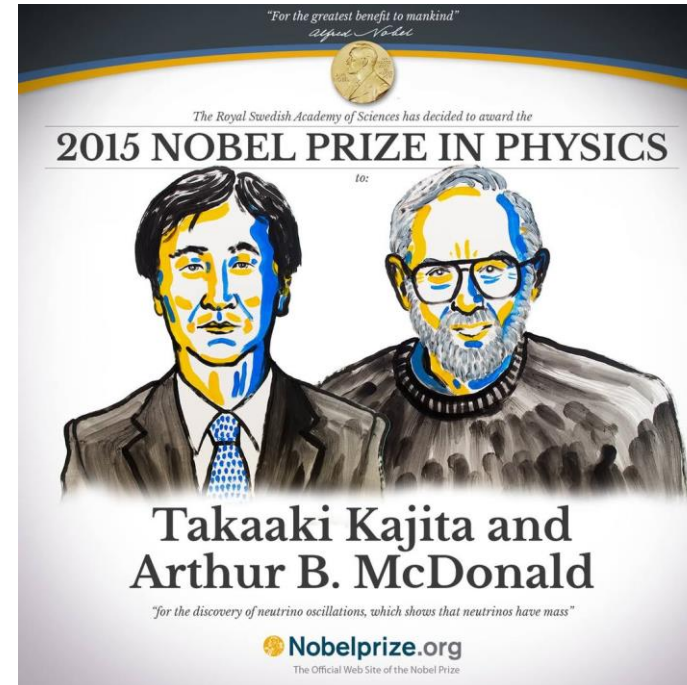


# Discovery of neutrino oscillations



The standard model of particle physics was built with:

- Minimal lepton sector
- No right-handed neutrino
- (Accidental) lepton number conservation
- Massless neutrinos



“for the discovery of neutrino oscillations,  
which shows that neutrinos have mass”

# The seesaw mechanism

$$\mathcal{L} \supset -m_D \bar{\nu}_D \nu_D - \frac{1}{2} m_R \overline{\nu_M^c} \nu_M + h.c.$$

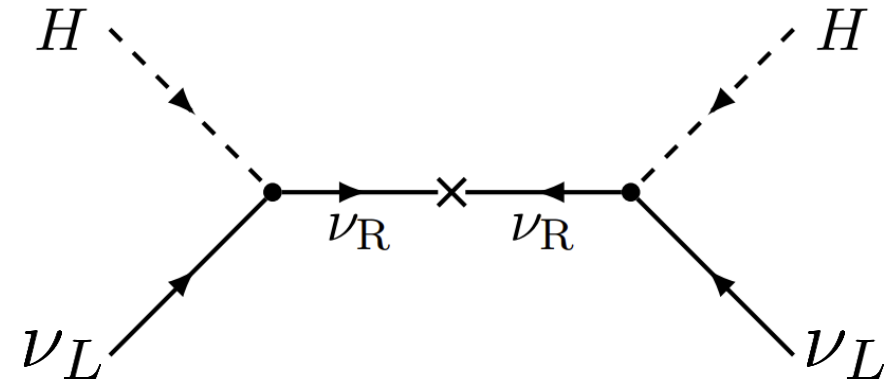
$$\nu_D = \nu_L + \nu_R, \quad \nu_M = \nu_R + \nu_R^c$$

Mass matrix for  $\nu_L, \nu_R$

$$\mathcal{M}_\nu = \begin{pmatrix} 0 & m_D \\ m_D & m_R \end{pmatrix}$$

$$\downarrow m_D \ll m_R$$

$$m_l \simeq \frac{m_D^2}{m_R}, \quad m_h \simeq m_R$$



$\nu_L$  : Left-handed neutrino  
 $H$  : Higgs  
 $\nu_R$  : Right-handed neutrino



# Vacuum neutrino masses as an explanation of oscillation data

$$H = \sqrt{p^2 + m_\nu^2} \approx p + \frac{m_\nu^2}{2E_\nu}$$

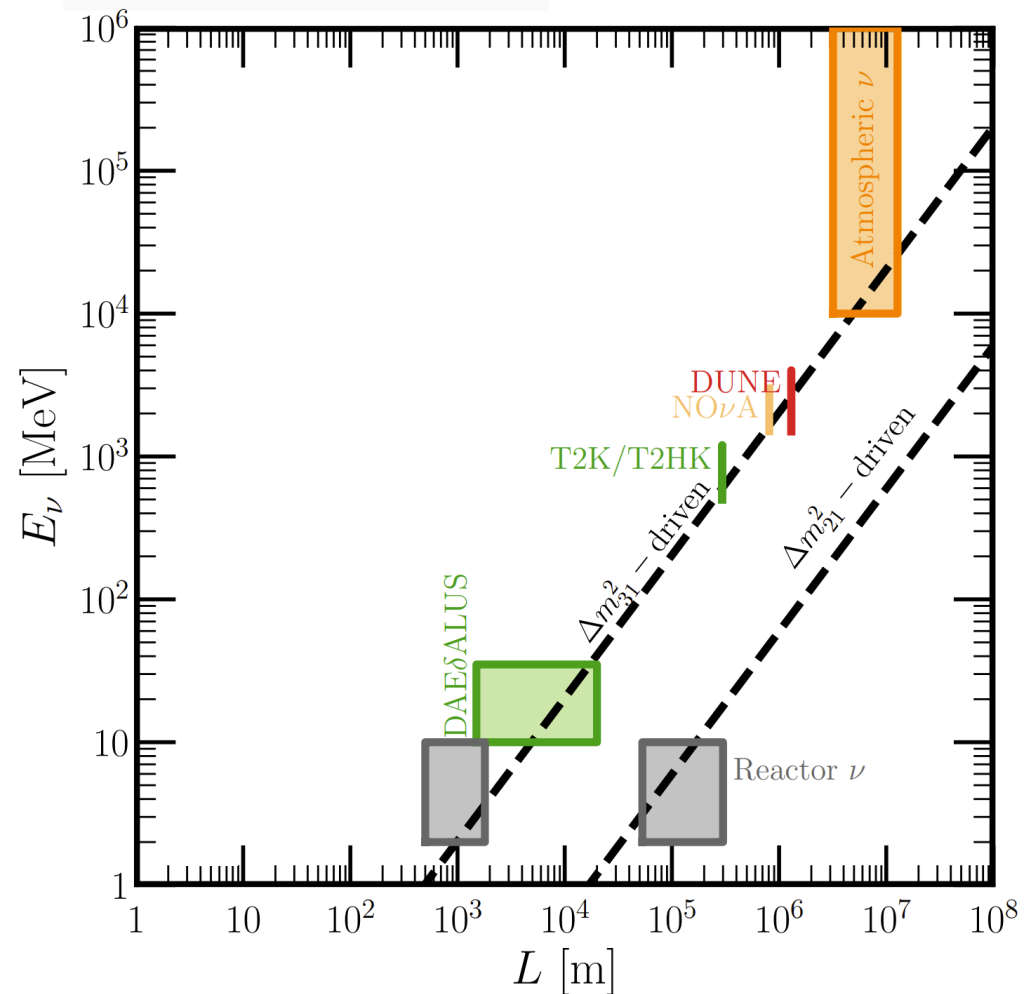
$$|\nu_i, t\rangle = e^{-iHt} |\nu_i\rangle$$

$$\text{Oscillation phase} \propto \frac{\Delta m_{ij}^2 L}{E_\nu}$$

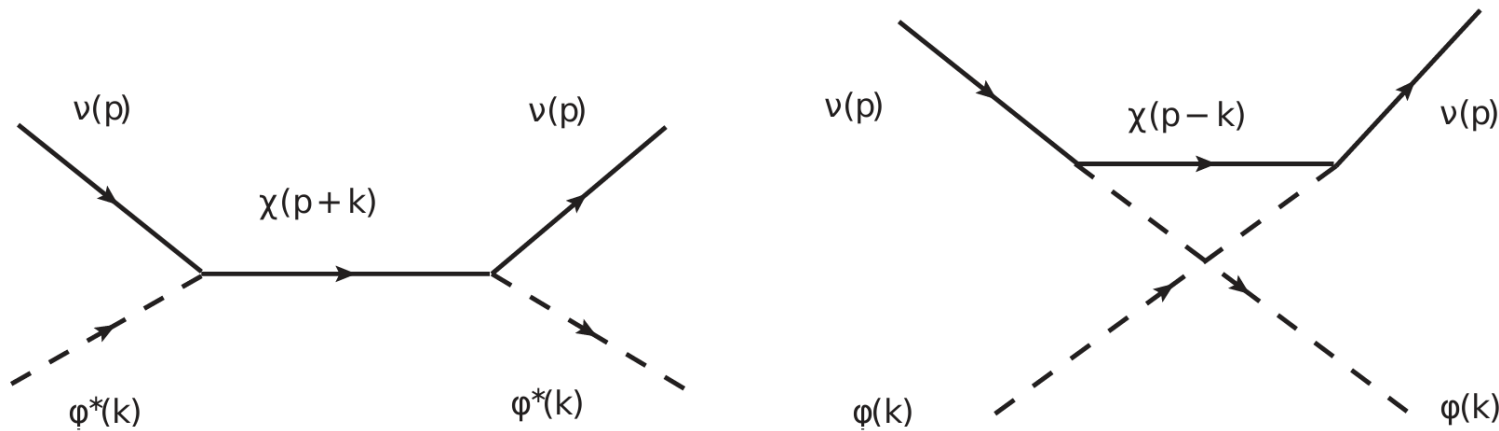


What if  $m_\nu = 0$  but with

$$H = p + V, \quad V = \frac{m_\nu^{\text{eff}}}{2E_\nu} ?$$



# A specific realization of “dark” neutrino mass



$\nu$  : Neutrinos  
 $\chi$  : Fermionic mediators  
 $\varphi$  : Scalar dark matter  
 $\nu, \chi$  have zero bare mass

Cold gas of dark matter particles:  
(Forward scattering)

$$m_\nu^2 \propto \frac{\rho_\phi}{m_\phi^2} \frac{y(y - \epsilon)}{y^2 - 1}, \quad y = \frac{2E_\nu m_\phi}{m_\chi^2}, \quad \epsilon = \frac{n_\phi - \bar{n}_\phi}{n_\phi + \bar{n}_\phi}$$

Classical scalar field background:

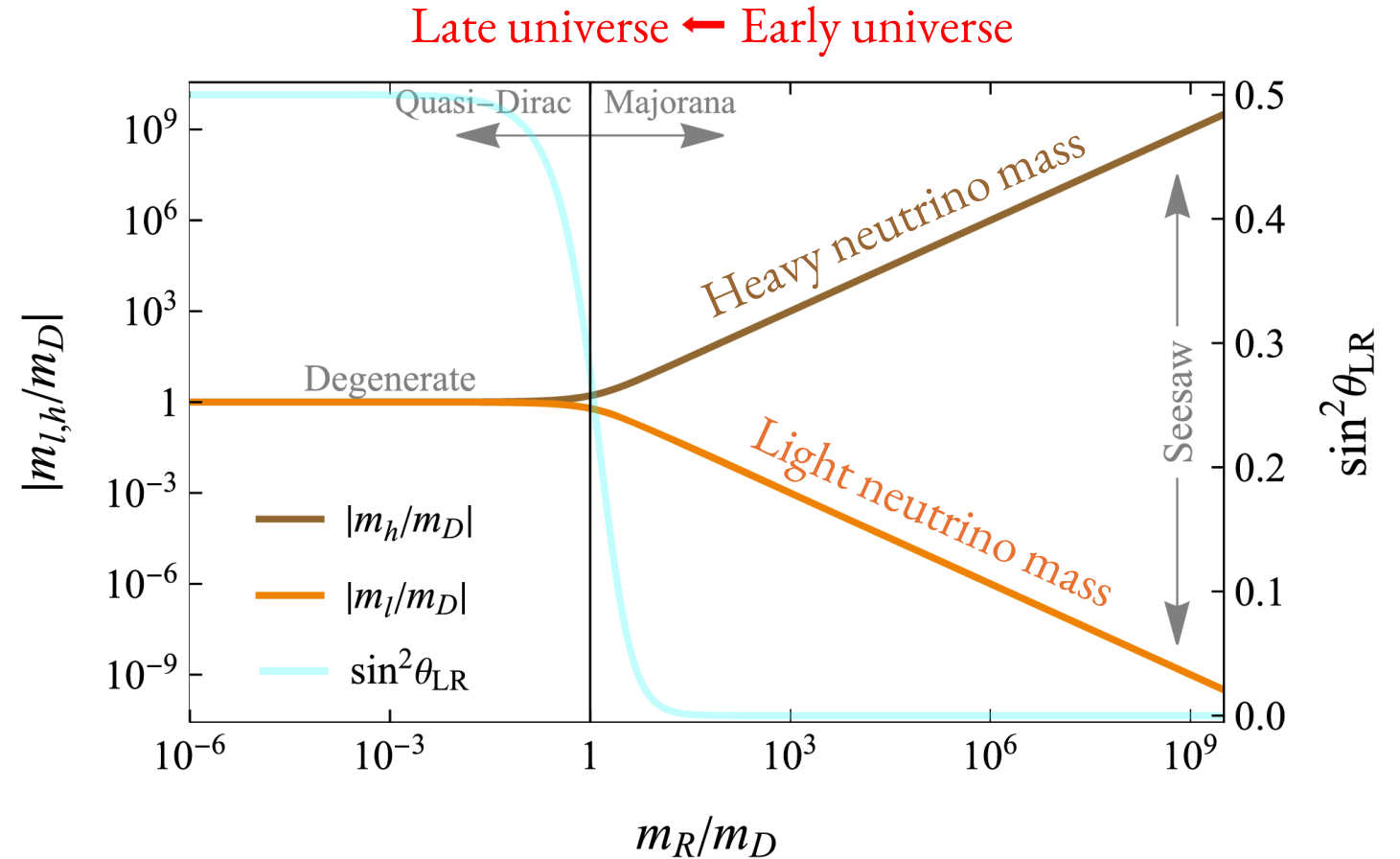
$$m_\nu^2 \propto \frac{\rho_\phi}{m_\phi^2} \cos^2(m_\phi t) \quad (\text{Relevant to ultralight dark matter})$$

# Another realization of “dark” neutrino mass

$$\mathcal{L} \supset -m_D \bar{\nu}_D \nu_D - \frac{1}{2} m_R \overline{\nu_M^c} \nu_M + h.c.$$

If the Majorana mass is due to couplings to dark matter:

$$m_R = g\phi_0 \cos(m_\phi t)$$



$$\begin{pmatrix} \nu_l \\ \nu_h \end{pmatrix} = \begin{pmatrix} \cos \theta_{LR} & \sin \theta_{LR} \\ -\sin \theta_{LR} & \cos \theta_{LR} \end{pmatrix} \begin{pmatrix} \nu_L^c \\ \nu_R \end{pmatrix} + \text{H.c.}$$

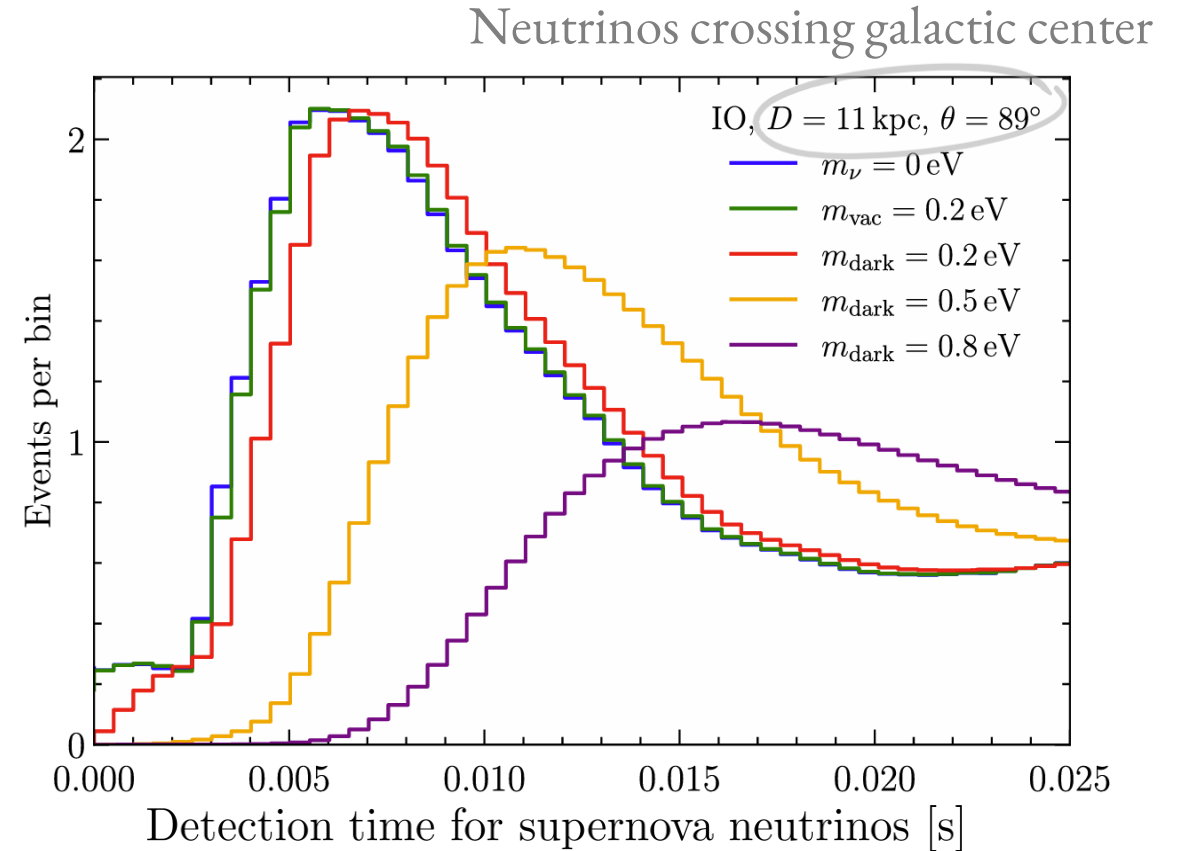


# Testing the mass origin with supernova neutrinos

Rate for galactic core-collapse supernovae is low,  $\sim O(1)$ /century Adams et al. (2013)

Arrival time delay effect is pronounced for:

- Large “dark” mass
- Supernovae near galactic center (even lower rate)



# Tests with oscillation experiments

Aiming for the entire ultralight (wavelike) mass range  $10^{-19}\text{eV} \lesssim m_\phi \ll 10\text{eV}$

# Several time and length scales



Parametrization for mass-squared difference:  $\Delta m_{ij}^2 = \Delta m_{ijD}^2(\mathbf{x}) \cos^2(m_\phi t)$

Oscillation period:  $T_\phi = \frac{\pi}{m_\phi} = 5.7\text{hr} \left( \frac{10^{-19}\text{eV}}{m_\phi} \right)$

Typical duration of oscillation experiments:  $T_{\text{exp}} \gtrsim \mathcal{O}(10)\text{days}$

Time-averaged probabilities

Dark matter coherence length:  $\lambda_{\text{dB}} = 1.24\text{au} \left( \frac{10^{-14}\text{eV}}{m_\phi} \right) \left( \frac{200\text{km/s}}{v} \right)$

Crossing distance of Earth during an experiment:  $l_\oplus = 1.16\text{au} \left( \frac{T_{\text{exp}}}{10\text{days}} \right) \left( \frac{v_\oplus}{200\text{km/s}} \right)$

Space-averaged probabilities for  $m_\phi \gg 10^{-14}\text{eV}$



# Strategy

For  $m_\phi \ll 10^{-14}\text{eV}$ , DM field has constant amplitude during an experiment:

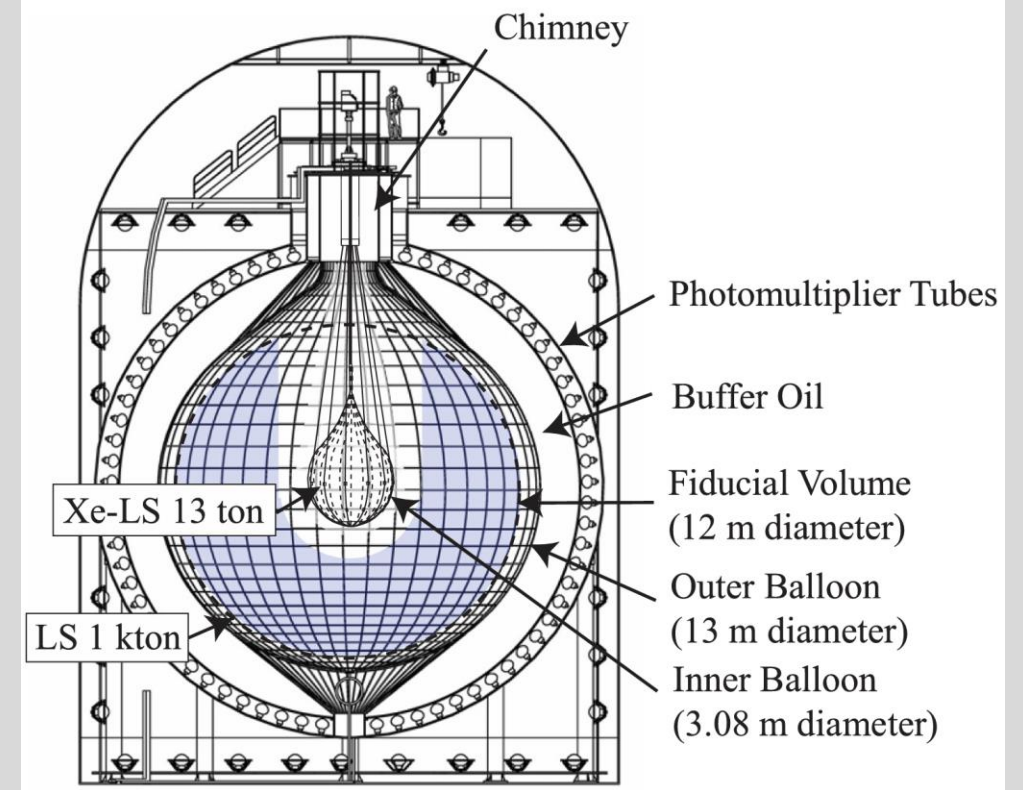
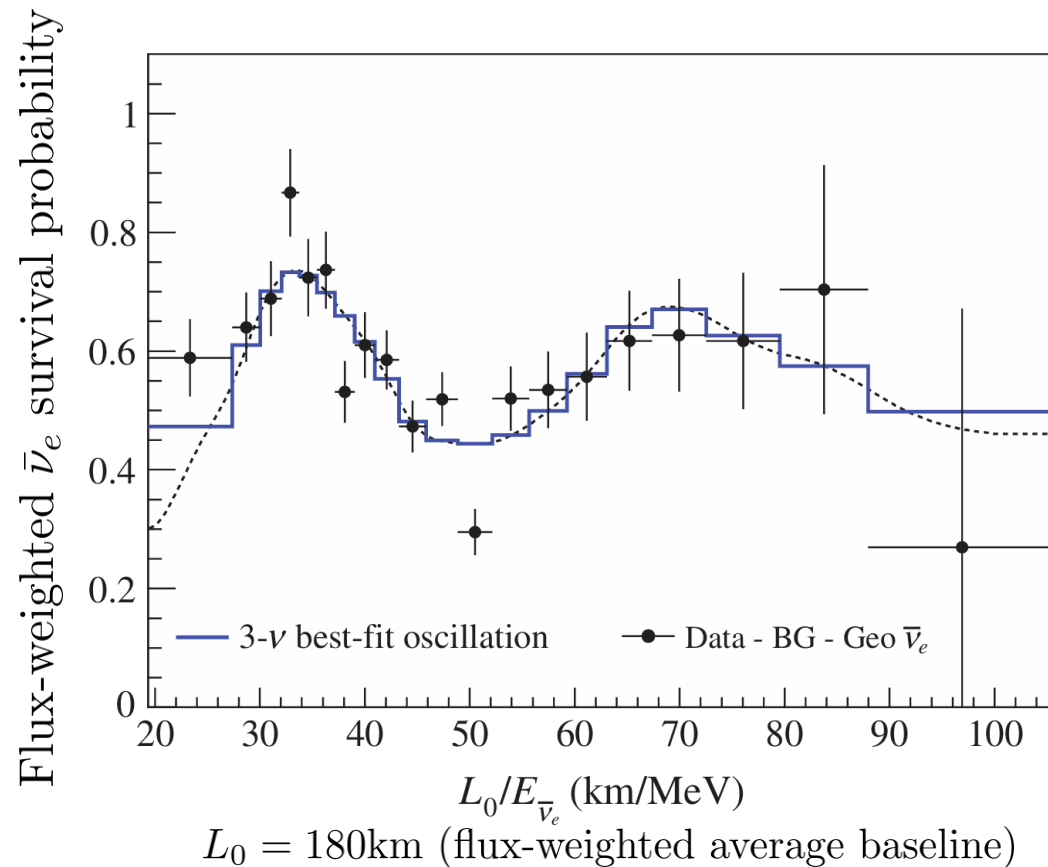


1. Find time-averaged formulae for flavor oscillations
2. Fit the formula with oscillation data (e.g., KamLAND)

For  $m_\phi \gg 10^{-14}\text{eV}$ , DM field has stochastic amplitudes in different de Broglie patches:

1. Model spatial fluctuations of ultralight dark matter
2. Take spatial average of the time-averaged formula
3. Compare the formula with oscillation data (if needed)

# Long-baseline reactor experiment: KamLAND



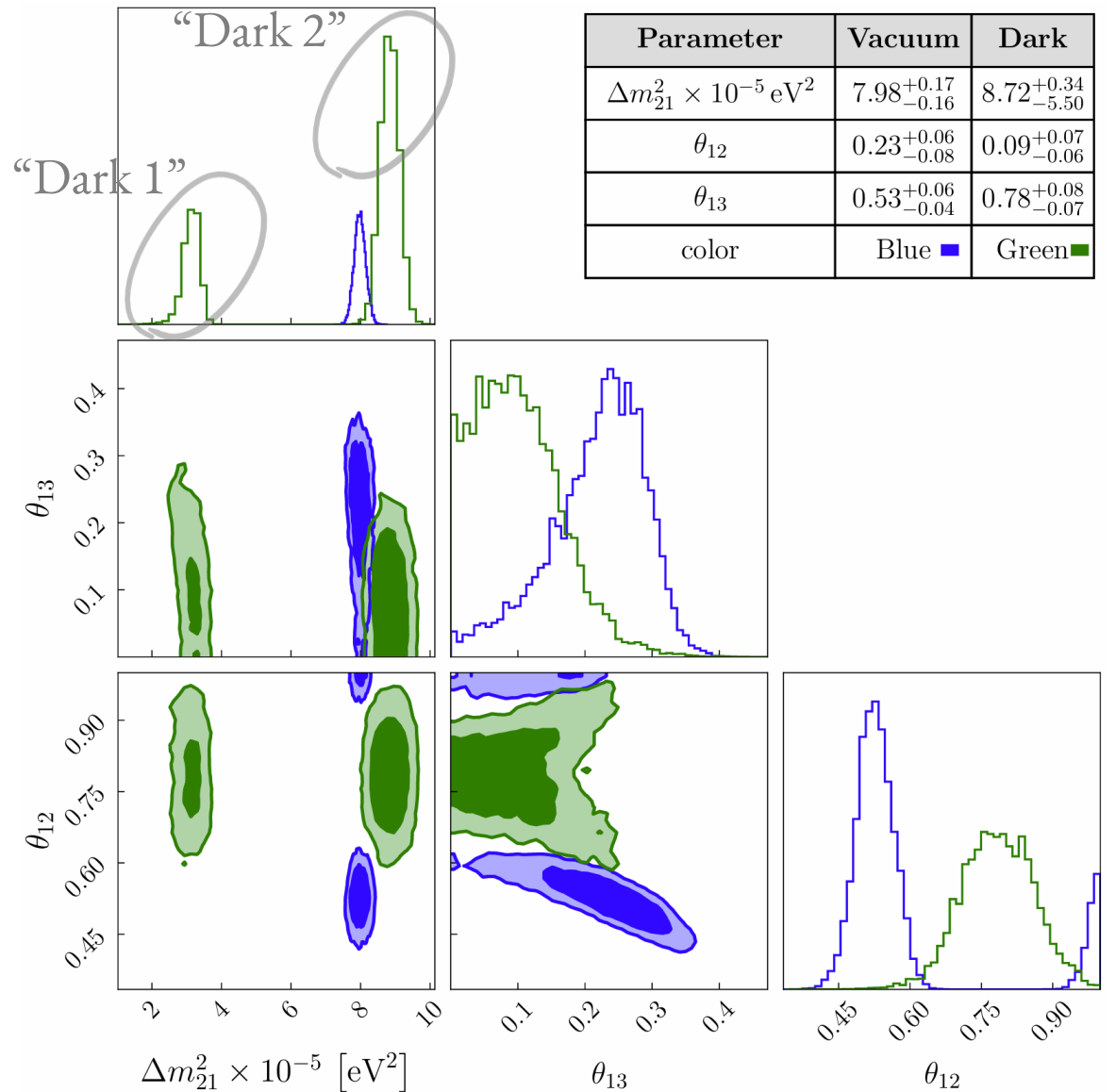
Located at 1km underground, Hida, Japan  
Detected antineutrinos from >50 reactors (before 2013)  
Sensitive to  $\Delta m_{21}^2$ ,  $\theta_{12}$ ,  $\theta_{13}$

# Chi-square analysis

Time-averaged survival probability:

$$P_{ee} = 1 - \frac{1}{2} \cos^4 \theta_{13} \sin^2(2\theta_{12}) [1 - J_0(X_{21D}) \cos X_{21D}] - \frac{1}{2} \sin^2(2\theta_{13}) [1 - J_0(X_{32D}) \cos X_{32D}]$$

$$X_{ijD} = \frac{\Delta m_{ij}^2 L}{4E_\nu}, \quad J_0(x) \text{ is Bessel function}$$

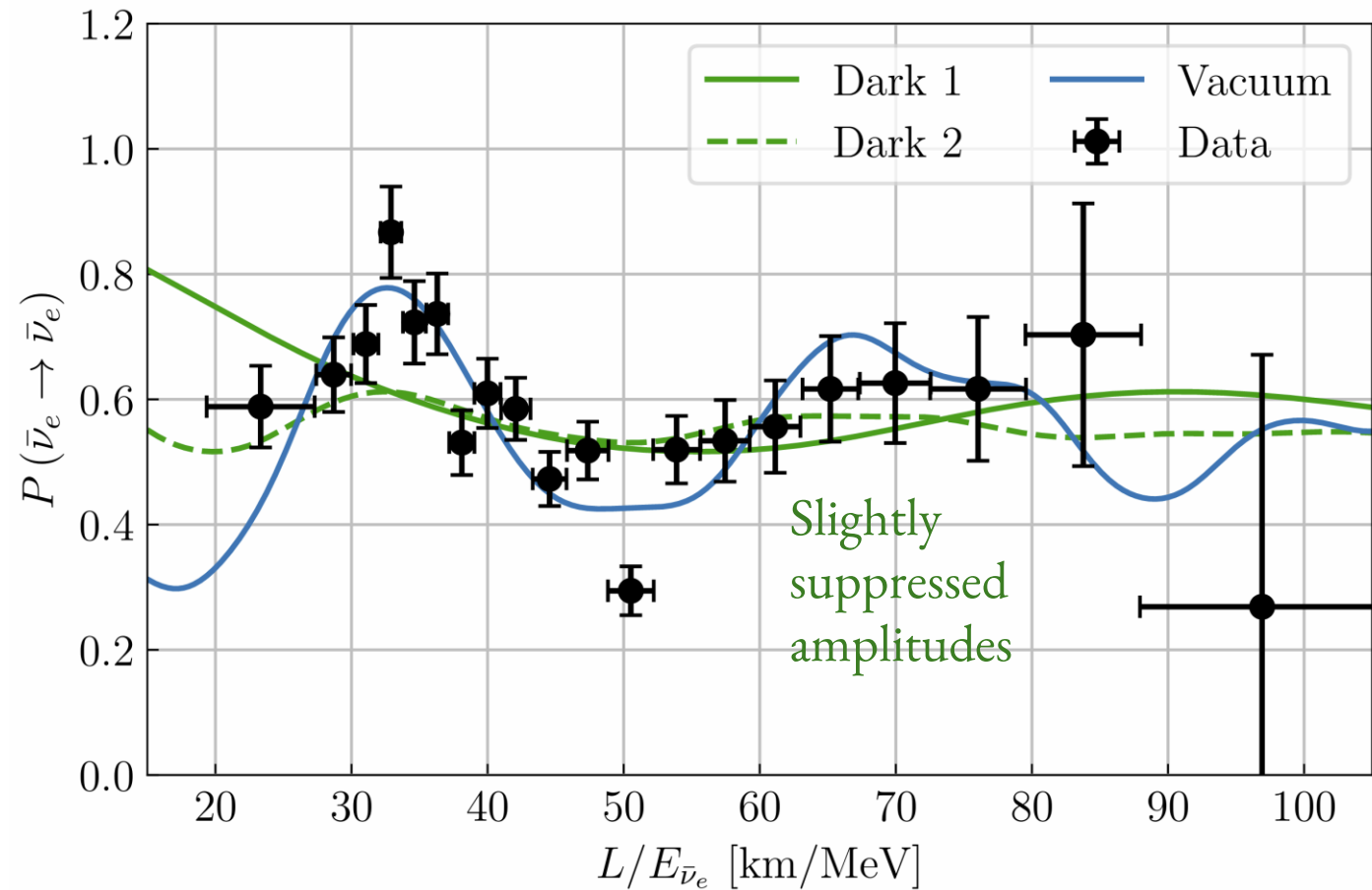


# Survival probabilities with best-fit parameters

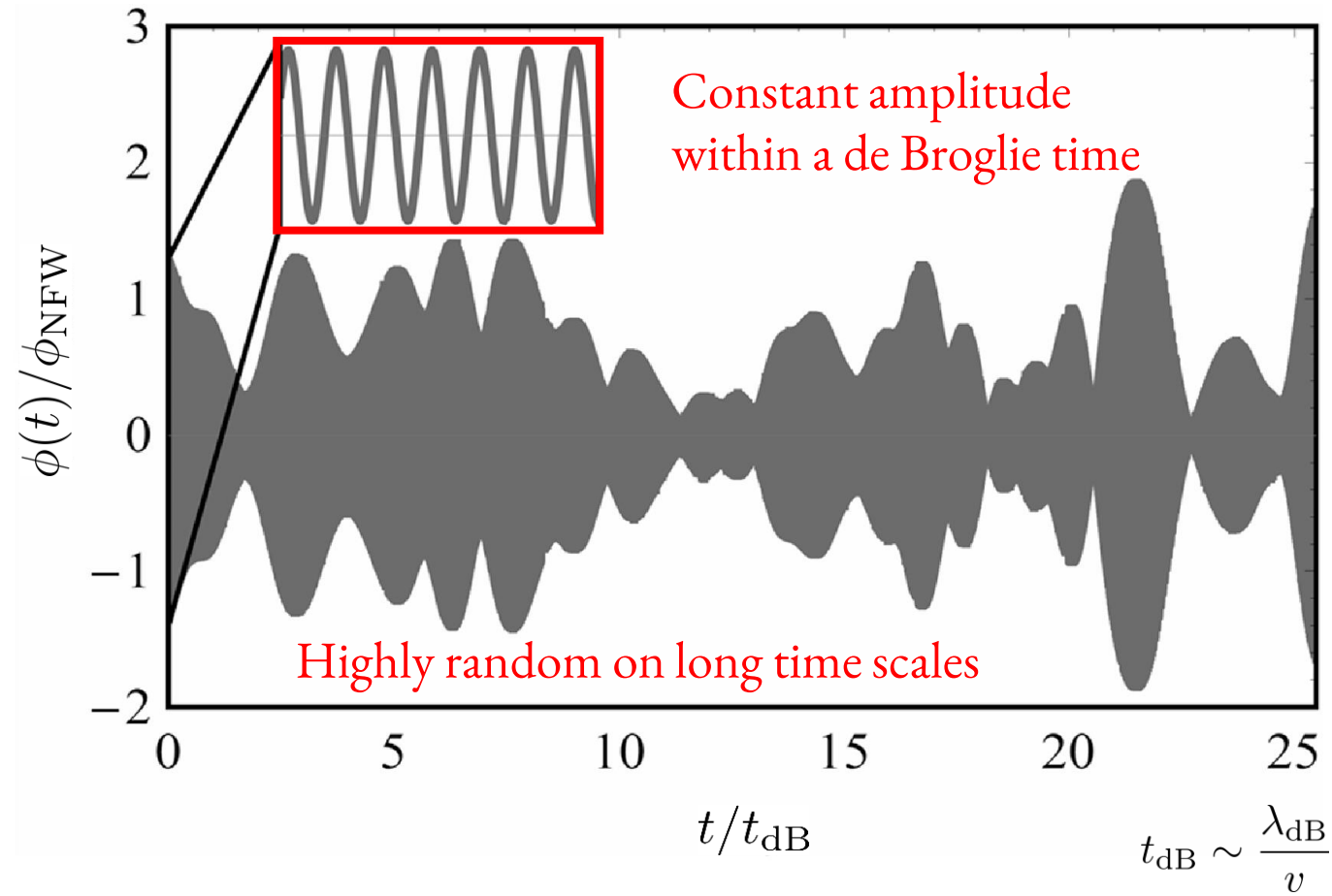
$$\chi^2_{\min, \text{vac}} = 35.5$$

$$\chi^2_{\min, \text{dark}} = 61.7$$

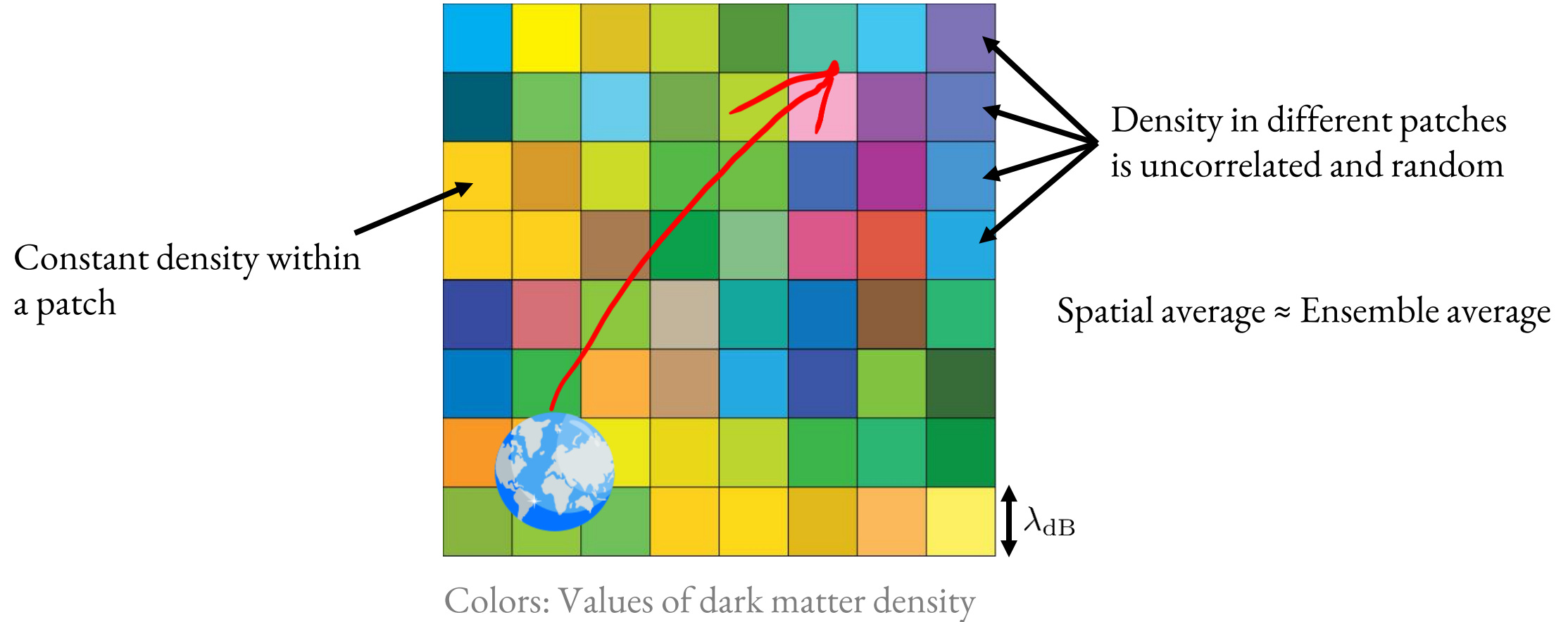
Vacuum (constant) mass  
is favored at  $4.5\sigma$



# Stochastic fluctuations for ultralight dark matter



# Earth crossing through different de Broglie patches



# Suppressed oscillation behaviors

$$P_{ee} = 1 - \cos^4 \theta_{13} \sin^2(2\theta_{12}) F_{21} - \sin^2(2\theta_{13}) F_{32}$$

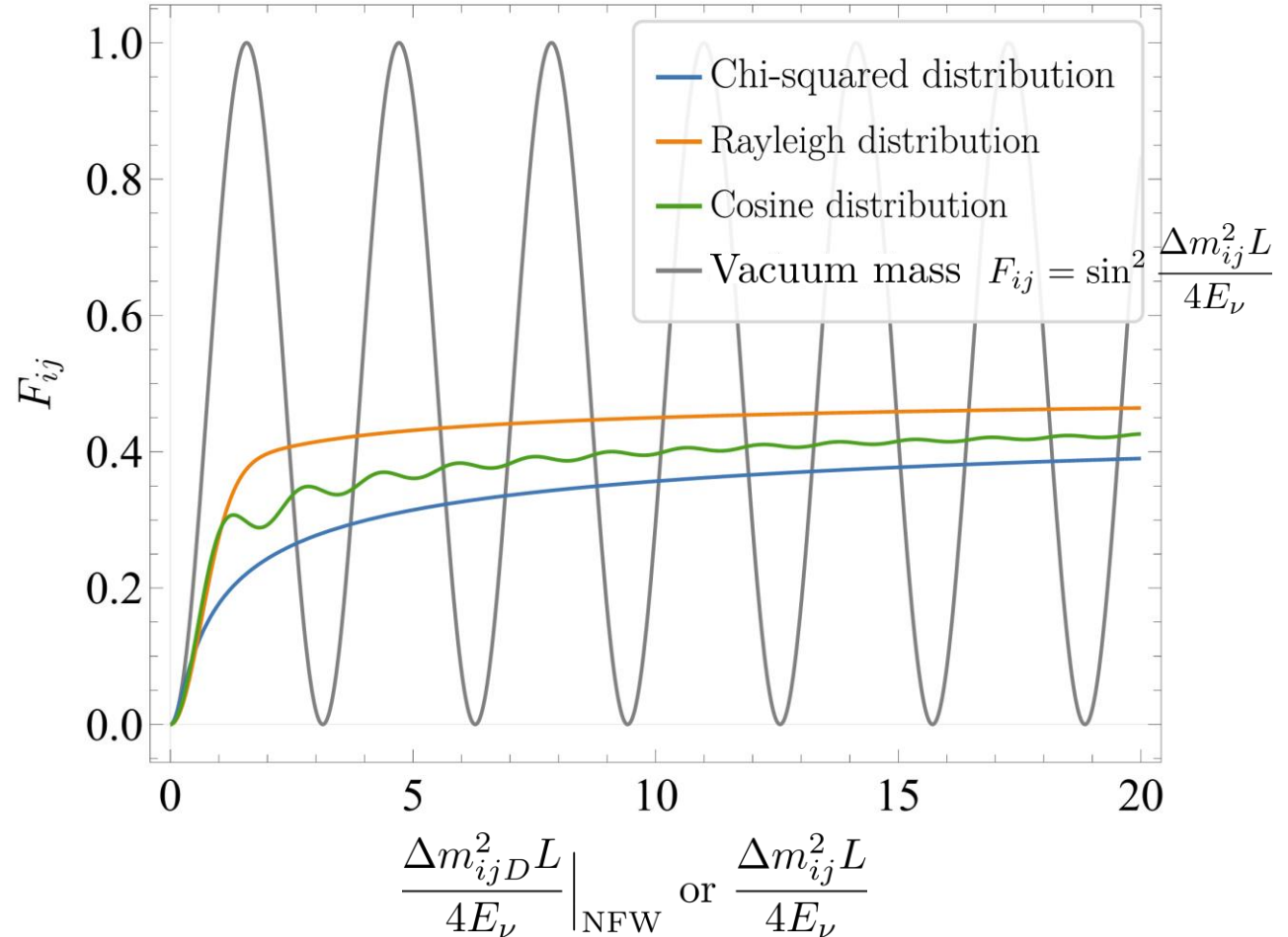
Time-averaged component:

$$F_{ij} = \frac{1}{2} \left[ 1 - J_0 \left( \frac{\Delta m_{ijD}^2 L}{4E_\nu} \right) \cos \left( \frac{\Delta m_{ijD}^2 L}{4E_\nu} \right) \right]$$

Distributions of mass-squared differences in different de Broglie patches:

$$\frac{\Delta m_{ijD}^2}{\Delta m_{ijD}^2|_{\text{NFW}}} \sim \chi^2(1), \text{ Rayleigh, } 1 - \cos(k_\theta \theta)$$

Flavor oscillations in terms of distances are suppressed in a model-independent way!





# Conclusions



- For  $10^{-19} \lesssim m_\phi \ll 10^{-14} \text{eV}$ ,

KamLAND disfavors “dark” neutrino mass by more than  $4\sigma$ .

- For  $10^{-14} \text{eV} \ll m_\phi \ll 10 \text{eV}$ ,

Stochastic DM fluctuations suppress neutrino oscillations.

- Ultralight/wavelike dark matter is unlikely to account for neutrino mass.

